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Causality Constraints in Statistical Hadron Production

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QCD Hadronization and the Statistical Model

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Two Questions:

- Why is strangeness **suppressed** for hadron production in **elementary collisions**?
- Why is strangeness **not suppressed** for hadron production in **nuclear collisions**?

One Answer:

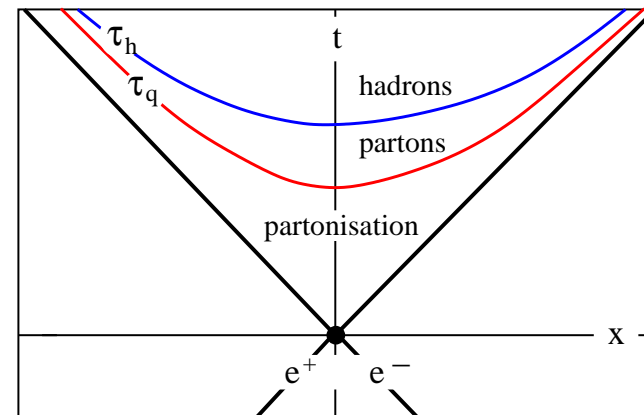
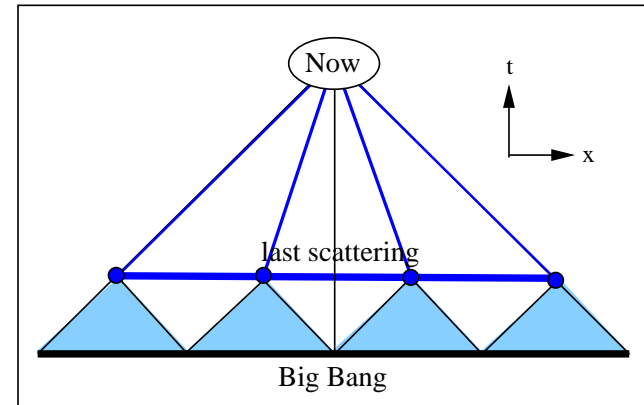
- **Causality constraints** restrict production to separate small spatial regions, and within a given region
 - there is **one pair** of strange quarks (elementary collisions)
 - there are **many pairs** of strange quarks (nuclear collisions)

Joint work with Paolo Castorina

Recall horizon problem of cosmology:

cosmic background radiation
is today 2.72548 ± 0.00057 °K
up to fourth decimal
throughout observable universe
why? It comes from
causally disconnected regions

\exists horizon problem also
in high energy hadroproduction;
inside-outside cascade
locally after time τ_q :
partonisation (equilibrium QGP)
locally after time τ_h :
hadronisation

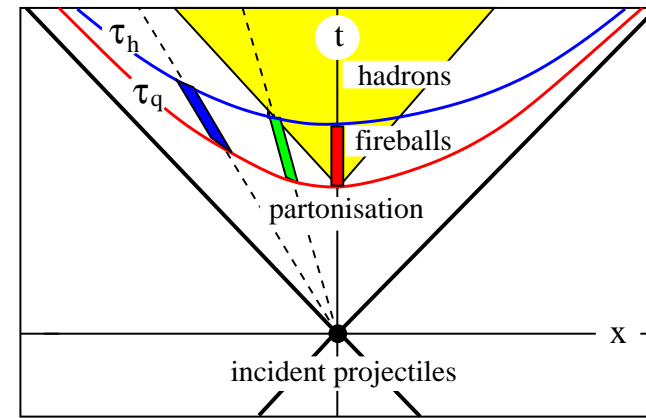


bubble of partonic medium of proper time τ
 with $\tau_q < \tau < \tau_h$: **fireball**;
 fireballs at different spatial
 rapidities η

$$t = \tau \cosh \eta, \quad x = \tau \sinh \eta,$$

with transition lines

$$t^2 - x^2 = \tau^2$$



red fireball ($\eta = 0$) - causality region yellow

green fireball ($\eta = \eta_d$) - one common x-t point with red

blue fireball ($\eta > \eta_d$) - outside causality region of red

for $\eta > \eta_d$, with $\tanh \eta_d = (\tau_h^2 - \tau_q^2) / (\tau_h^2 + \tau_q^2)$

forward and backward fireballs are out of communication
 with central fireball

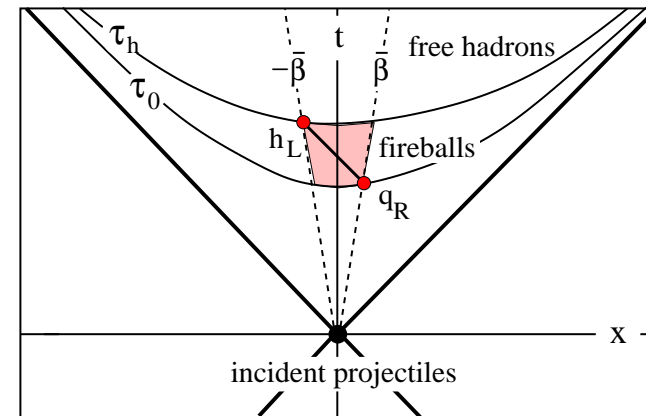
examples:

$$\tau_q = 1 \text{ fm}, \tau_h = 2 \text{ fm} \rightarrow \eta_d = 0.7$$

$$\tau_q = 1 \text{ fm}, \tau_h = 7 \text{ fm} \rightarrow \eta_d = 2$$

at RHIC and LHC, hadronisation occurs through causally disjoint fireballs

so far, have neglected spatial size:
 what is the size of a fireball?
 define through causal connectivity
 require: the most separate points
 can still communicate



spatial diameter d of fireball in cms at hadronisation time

$$d = \sqrt{\frac{\tau_h}{\tau_q}} (\tau_h - \tau_q)$$

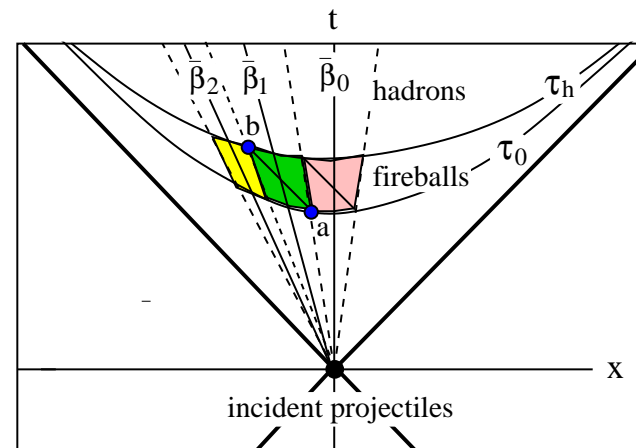
examples for different
hadronisation times:

τ_h [fm]	β	η	$r = d/2$ [fm]
2	0.33	0.35	0.7
3	0.50	0.55	1.7
4	0.60	0.69	3.0
5	0.67	0.81	4.5

denote average cms velocity of of central fireball by $\bar{\beta}_0 = 0$
can then partition production region into successive causally
disjoint fireballs, of velocities

$$\bar{\beta}_n = \frac{\tau_h^{2n} - \tau_q^{2n}}{\tau_h^{2n} + \tau_q^{2n}}$$

$$n = 0, 1, 2, \dots$$



statistical hadronisation:

- massive colorless clusters distributed over rapidities,
each decays statistically
- mass and charge distributions of clusters again statistically
 \Rightarrow equivalent global cluster
- $V = \Sigma V_i, Q = \Sigma Q_i$; large enough for thermodynamics

$$\langle n_j \rangle_{\text{primary}} = \frac{VTm_j^2}{2\pi^2} \gamma_s^{n_j} K_2 \left(\frac{m_j}{T} \right)$$

with $\gamma_s \sim 0.5 - 0.7$ to account for strangeness suppression
in elementary collisions

Causality constraints \Rightarrow equivalent global cluster **not allowed**
in elementary collisions

\sim one pair of strange particles/causality cluster, strangeness
must be conserved locally

Local strangeness suppression [Hamieh, Redlich, Tounsi 2000]

- canonical strangeness conservation

$$\frac{Z_{\text{can}}(T, V, S)}{Z_{\text{gcan}}(T, V, \langle S \rangle)} < 1$$

- local strangeness conservation in a strangeness correlation
volume $V_c \ll V$

\Rightarrow 3 parameters: T, V, V_c

accounts for data, with V_c in place of γ_s

Here: $V_c =$ volume of causal correlation cluster

data analysis: $d_c \sim 2 \text{ fm} \rightarrow \tau_h \sim 2 - 3 \text{ fm}$

- causality constraints provide basis for
local strangeness conservation

Nuclear collisions:

space-time overlap of many causally connected clusters at each $\bar{\beta}_i$ provides equivalent global cluster containing many $s\bar{s}$ pairs

moreover, superposition increases parton density & hence τ_h :
larger correlation volume

for RHIC & LHC, \sim no strangeness suppression

So far: strangeness correlation on hadronic level, at τ_h .

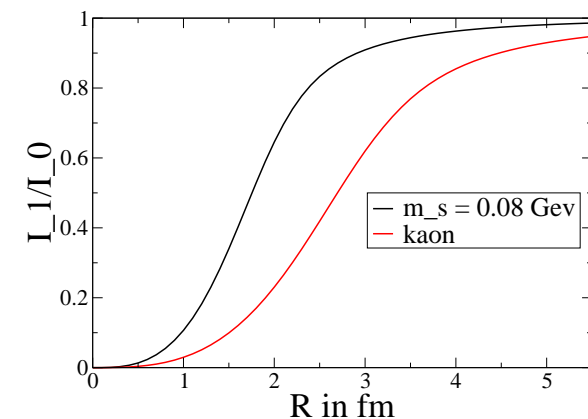
what about $\phi(s\bar{s})$?

possible solution:

local strangeness correlation

on quark level, prior to hadronisation

suppression factor for $m_s = 80$ MeV



Conclusions

- Onset of thermalisation implies the existence of causally disconnected space-time formation regions
- For rare species (strangeness), this requires local conservation in elementary collisions, rules out equivalent global cluster
- In nuclear collisions, superposition of causally connected clusters (at the same space-time location) allows global clusters, removes strangeness suppression
- open horizon problem remains: CGC \sim false vacuum,
glasma \sim inflation?