# Extracting information on nuclear charge distributions from isotope shift studies

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- Michel Godefroid and Cedric Nazé, Université Libre de Bruxelles
- Gediminas Gaigalas, Vilnius University
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- Isotope shifts
- $\bullet$  Reformulated field shift and implementation in  ${\rm GRASP2K}$
- Applications
  - Refined extraction of  $\delta \langle r^2 \rangle$  values
  - Higher order corrections from realistic charge distrbutions
  - Simultaneous extraction of  $\delta\langle r^2
    angle$  and  $\delta\langle r^4
    angle$  values
- Conclusions and Outlook



#### Isotope shifts



$$\delta \nu_{k,\text{IS}}^{A,A'} = \delta E_u^{A,A'} - \delta E_l^{A,A'}$$

$$= \underbrace{\Delta K_{k,\text{MS}}}_{\text{MM'}} \underbrace{\frac{M' - M}{MM'}}_{\text{Mass Shift}} + \underbrace{\sum_{N} F_{k,N} \delta \langle r^{2N} \rangle^{A,A'}}_{\text{Field Shift}}$$

$$\sum_{N} F_{k,N} \delta \langle r^{2N} \rangle^{A,A'} = \delta \nu_{k,\mathrm{IS}}^{A,A'} - \Delta K_{k,\mathrm{MS}} \frac{M'-M}{MM'} = \delta \nu_{k,\mathrm{FS}}^{A,A'}$$

The possibility to extract accurate  $\delta \langle r^{2N} \rangle^{A,A'}$  ultimately depends on how well we measure experimental  $\delta \nu_{k,\text{IS}}^{A,A'}$  and calculate  $F_{k,N}, \Delta K_{k,\text{MS}}$ .



#### Reformulated field shift and implementation in $\operatorname{GRASP2K}$

J. Ekman, P. Jönsson, M. Godefroid, C. Nazé and G. Gaigalas, to be submitted to CPC



The first order level field shift is given by

$$\delta E_{i,\mathrm{FS}}^{(1)\mathcal{A},\mathcal{A}'} = -\int_{\mathbf{R}^3} \left[ V_{\mathcal{N}}^{\mathcal{A}}(\mathbf{r}) - V_{\mathcal{N}}^{\mathcal{A}'}(\mathbf{r}) \right] \rho_i^{\mathbf{e}}(\mathbf{r}) d^3\mathbf{r},$$

where

- $V_A(\mathbf{r})$  and  $V_{A'}(\mathbf{r})$  are the electronic potentials due to (different) charge distributions in the two isotopes A and A'.
- $\rho_i^e(\mathbf{r})$  is the electron density distribution within the nuclear volume



Assuming a homogenous charge distribution (hard sphere) it can be shown that the electron density distribution for level i inside the nuclear volume to a very good approximation can be expanded as [1]:

$$\rho_i^e(\mathbf{r}) \approx b_{i,0} + b_{i,2}r^2 + b_{i,4}r^4 + b_{i,6}r^6.$$

[1] Blundell et al., Journal of Physics B 20, 3663 (1987)

In fact, the polynomial expansion also holds for more realistic Fermi-type charge distributions:





Electron density for the  $1s^22s^2S$  level in Li-like <sup>142</sup>Nd ( $\approx$  7 % fall-off within nuclear volume)



#### Reformulated field shift

Using the polynomial expansion  $\rho_i^e(\mathbf{r}) \approx b_{i,0} + b_{i,2}r^2 + b_{i,4}r^4 + b_{i,6}r^6$  and following the work by Seltzer and Blundell *et al.* the first order frequency field shift can then be written

$$\delta\nu_{k,\rm FS}^{(1)A,A'} = F_{k,0}\delta\langle r^2\rangle^{A,A'} + F_{k,2}\delta\langle r^4\rangle^{A,A'} + F_{k,4}\delta\langle r^6\rangle^{A,A'} + F_{k,6}\delta\langle r^8\rangle^{A,A'}$$

where

$$F_{k,n} = \frac{4\pi}{(n+2)(n+3)} Z \frac{\Delta b_n}{h}, \quad \Delta b_n = b_{u,n} - b_{l,n}.$$

Seltzer, Phys. Rev. **188**, 1916 (1969) Blundell *et al.*, Journal of Physics B **20**, 3663 (1987)



## Routines for the derivation of the electronic factors $F_N$ are implemented in GRASP2K through the program RIS4.

Level J Parity Field shift electronic factors and average point discrepancy in fit

		FO (GHz/fm^2)	F2 (GHz/fm^4)	F4 (GHz/fm^6)	F6 (GHz/fm^8)	Disc. (per mille)
1	1/2 -	0.1155473544D+06	-0.9457001725D+02	0.3054728900D+00	-0.6051012375D-03	0.0020
1	3/2 -	0.1152739390D+06	-0.9435715260D+02	0.3047762374D+00	-0.6037193610D-03	0.0020

## J. Ekman, P. Jönsson, M. Godefroid, C. Nazé and G. Gaigalas, to be submitted to CPC



Decomposition of the level field shifts of the  ${}^2S_{1/2}$  and  ${}^2P_{1/2,3/2}$  states in Li-like  ${}^{150,142}$ Nd from a n = 5 + Breit calculations.

	$1s^2 2s \ ^2S_{1/2}$	1s <sup>2</sup> 2p <sup>2</sup> P <sub>1/2</sub>	1s <sup>2</sup> 2p <sup>2</sup> P <sub>3/2</sub>
$\mathcal{F}_0\delta\langle r^2\rangle^{150,142}$	656.78	614.67	613.23
$\mathcal{F}_2\delta\langle r^4 angle^{150,142}$	-33.10	-30.96	-30.89
$\mathcal{F}_4\delta\langle r^6 angle^{150,142}$	5.89	5.50	5.49
$\mathcal{F}_{6}\delta\langle r^{8} angle^{150,142}$	-0.66	-0.62	-0.62
$\delta E_{\rm FS}^{(1)150,142}$	628.91	588.60	587.21



Decomposition of the frequency field shifts in meV of the  ${}^2P_{1/2} - {}^2S_{1/2}$  and  ${}^2P_{3/2} - {}^2S_{1/2}$  transitions in Li-like  ${}^{150,142}$ Nd from a resulting n = 5 + Breit calculation.

	$^{2}P_{1/2} - ^{2}S_{1/2}$	$^{2}P_{3/2} - ^{2}S_{1/2}$
$F_0\delta\langle r^2\rangle^{150,142}$	-42.11	-43.55
$F_2\delta\langle r^4\rangle^{150,142}$	2.13	2.20
$F_4\delta\langle r^6 angle^{150,142}$	-0.38	-0.39
$F_6\delta\langle r^8 angle^{150,142}$	0.04	0.04
$\delta  u_{ m FS}^{(1)150,142}$	-40.31	-41.70



Comparison of different approaches of frequency field shifts calculations of the  ${}^2P_{1/2} - {}^2S_{1/2}$  transition in Li-like  ${}^{150,142}$ Nd as a function of active orbital set.

	Variational	Reformulated	1st. ord. pert.
DF	-40.44	-40.39	-40.39
<i>n</i> = 2	-40.43	-40.55	
<i>n</i> = 3	-40.44	-40.51	
<i>n</i> = 4	-40.45	-40.53	
<i>n</i> = 5	-40.45	-40.53	
n = 5 + Breit	-40.30	-40.31	



#### Effect from deformation in Li-like <sup>150,142</sup>Nd

#### From a n=5 + Breit calculation. Values in meV.

	$\delta \langle r^2 \rangle$ [fm <sup>2</sup> ]	$\delta \langle r^4 \rangle$ [fm <sup>4</sup> ]	$\delta \langle r^6 \rangle$ [fm <sup>6</sup> ]	$\delta \langle r^8 \rangle$ [fm <sup>8</sup> ]
<sup>150</sup> Nd (sph.)	1.29094	79.4628	4 375.68	248 818
$^{150}$ Nd ( $\beta_{20} = 0.2$	8) 1.29094	94.5064	6 016.55	386 461
<sup>2</sup> P	$S_{1/2} - S_{1/2}$ (sph.)	$^{2}P_{1/2} - ^{2}S$	$S_{1/2} \ (\beta_{20} = 0.28)$	) Difference
$F_0\delta\langle r^2\rangle^{150,142}$	-42.11	r.	-42.1	1 0.00
$F_2\delta\langle r^4\rangle^{150,142}$	2.13		2.5	4 0.40
$F_4\delta\langle r^6\rangle^{150,142}$	-0.38		-0.5	2 -0.14
$F_6\delta\langle r^8\rangle^{150,142}$	0.04		0.0	7 0.02
$\delta \nu_{\rm FS}^{(1)150,142}$	-40.31		-40.0	2 0.29

Effect from deformation: +0.29 meV, in good agreement with Kozhedub et al., Phys. Rev. A. **77**, 032501 (2008) Zubova et al., Phys. Rev. A. **90**, 062512 (2014)



## Applications: Refined extraction of $\delta \langle r^2 \rangle$ values



Rearranging, the reformulated field shift

$$\delta\nu_{k,\mathrm{FS}}^{(1)A,A'} = F_{k,0}\delta\langle r^2\rangle^{A,A'} + F_{k,2}\delta\langle r^4\rangle^{A,A'} + F_{k,4}\delta\langle r^6\rangle^{A,A'} + F_{k,6}\delta\langle r^8\rangle^{A,A'}$$

can be written as

$$\delta\nu_{k,\mathrm{FS}}^{(1)A,A'} = F_{k,0}\lambda_k^{A,A'}$$

where the so-called nuclear parameter  $\lambda_k^{\mathcal{A},\mathcal{A}'}$  is given by

$$\lambda_{k}^{\mathcal{A},\mathcal{A}'} = \delta \langle r^{2} \rangle^{\mathcal{A},\mathcal{A}'} + \frac{F_{k,2}}{F_{k,0}} \delta \langle r^{4} \rangle^{\mathcal{A},\mathcal{A}'} + \frac{F_{k,4}}{F_{k,0}} \delta \langle r^{6} \rangle^{\mathcal{A},\mathcal{A}'} + \frac{F_{k,6}}{F_{k,0}} \delta \langle r^{8} \rangle^{\mathcal{A},\mathcal{A}'}$$



In the following  $\delta \langle r^2 \rangle$  values are extracted from the nuclear parameter

$$\lambda_{k}^{A,A'} = \delta \langle r^{2} \rangle^{A,A'} + \frac{F_{k,2}}{F_{k,0}} \delta \langle r^{4} \rangle^{A,A'} + \frac{F_{k,4}}{F_{k,0}} \delta \langle r^{6} \rangle^{A,A'} + \frac{F_{k,6}}{F_{k,0}} \delta \langle r^{8} \rangle^{A,A'}$$

for the  $^2P_{1/2}-^2S_{1/2}$  transition in Li-like systemts using different Fermi charge distribution parameters for the two isotopes A and A'

$$ho(r, heta)=rac{1}{1+\exp(rac{r-c( heta)}{a})},\qquad c( heta)=c_0\left[1+eta_{20}Y_{20}( heta)
ight]$$

The skin diffusiness is often characterised by the parameter  $t = 4\ln(3)a$ .

25 August 2015



Fermi model charge distributions with  $\langle r^2 \rangle^{1/2} = 5.0$  fm and different skin diffusiness parameters t.



19 / 45

#### Constant density approximation: $\delta \langle r^2 \rangle = \lambda$





#### Spherical Fermi distributions with t = 2.3 fm for both isotopes





#### Homogenous distributions (t = 0.0 fm) for both isotopes





#### Spherical Fermi distributions with $t_A = 2.6$ and $t_{A'} = 2.3$ fm





#### Fermi distributions with $\beta_{2,A} = 0.3$ and $\beta_{2,A'} = 0.0$









Applications: Higher order corrections from realistic charge distrbutions

A. Papoulia, Master thesis, Lund University, 2015 B.G Carlsson, J. Ekman and A. Papoulia, in preparation



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Frequency field shifts are obtained using nuclear radial moments from

- Nuclear Hartree-Fock-Bogoliubov calculations (HFB)
- Two-parameter Fermi model with t = 2.3 fm (Fermi)
- a "correction" term is obtained as

$$\delta\nu_{k,\text{corr}}^{(1)A,A'} = \delta\nu_{k,\text{HFB}}^{(1)A,A'} - \delta\nu_{k,\text{Fermi}}^{(1)A,A'}$$
  
where  $\delta\langle r^2 \rangle_{\text{Fermi}}^{A,A'} = \delta\langle r^2 \rangle_{\text{HFB}}^{A,A'}$ 







The nuclear Hartree-Fock-Bogoliubov calculations (HFB) are perfomed using the Skyrme interaction with the UDF1 parameter set and the following software:

- HOSPHE (v2.00) for spherical nuclei B.G. Carlsson et al., Comput. Phys. Commun., 181, 1641 (2010)
- HFBTHO (2.00d) for axially symmetric deformed nuclei M.V. Stoitsov et al., Comput. Phys. Commun., 184, 1592 (2013)





Higher order corrections,  $\delta \nu_{k,\text{corr}}^{(1)A,142}$ , from realistic charge distributions for the  ${}^2P_{1/2} - {}^2S_{1/2}$  transition in Li-like Nd.





Higher order corrections,  $\delta \nu_{k,\text{corr}}^{(1)A,142}$ , from realistic charge distributions for the  ${}^2P_{1/2} - {}^2S_{1/2}$  transition in Li-like Nd.





Decomposition of expansion and corrections terms for the  ${}^{2}P_{1/2} - {}^{2}S_{1/2}$ transition in Li-like  ${}^{150,142}$ Nd. Experimental (statistical) uncertainty from Brandau et al., PRL **100** (2008) 973201



Higher order corrections,  $\delta \nu_{k,\text{corr}}^{(1)A,138}$ , from realistic charge distributions for selected transitions in neutral barium (Ba I).

Wave functions from C. Nazé, J.G. Li, and M. Godefroid, Phys. Rev. A **91**, 032511 (2015)





Higher order corrections,  $\delta \nu_{k,\text{corr}}^{(1)A,138}$ , from realistic charge distributions for selected transitions in neutral barium (Ba I).

C. Naze, J.G. Li, and M. Godefroid, Phys. Rev. A **91**, 032511 (2015)



#### Applications: Simultaneous extraction of $\delta \langle r^2 \rangle$ and $\delta \langle r^4 \rangle$ values

A. Papoulia, Master thesis, Lund University, 2015 B.G Carlsson, J. Ekman and A. Papoulia, in preparation



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Using the reformulated fiels shift and experimental isotope shifts for at least four analogue transitions in an isotope pair it is principle possible to extract  $\delta\langle r^2 \rangle$ ,  $\delta\langle r^4 \rangle$ ,  $\delta\langle r^6 \rangle$  and  $\delta\langle r^8 \rangle$ 

However, this is an extremely challenging task and would rely on experimental data and calculations without uncertainties.

Extracting  $\delta \langle r^2 \rangle$  and  $\delta \langle r^4 \rangle$  is more realistic, although extremely challenging, and would provide additional information of nuclear charge distributions.



25 August 2015

#### Simultaneous extraction of $\delta \langle r^2 \rangle$ and $\delta \langle r^4 \rangle$ values



The ratio  $\delta \langle r^4 \rangle^{A,A'} / \delta \langle r^2 \rangle^{A,A'}$  as a function of the deformation parameter  $\beta_2$  of the larger isotope.

 $\delta \langle r^2 \rangle^{1/2} = 5.0$  and 5.1 fm, respectively. t = 2.3 fm for both isotopes.

Approach 1:

Neglect  $\delta \langle r^6 \rangle$  and  $\delta \langle r^8 \rangle$  terms in the expression for the reformulated field shift:

$$\delta\nu_{k,\rm FS}^{(1)A,A'} \approx F_{k,0}\delta\langle r^2\rangle^{A,A'} + F_{k,2}\delta\langle r^4\rangle^{A,A'}$$

Using experimental isotope shifts for two transitions, 1 and 2, in an isotope pair,  $\delta \langle r^4 \rangle^{A,A'}$  is then given by

$$\delta \langle r^{4} \rangle^{A,A'} = \frac{\lambda_{1}^{A,A'} - \lambda_{2}^{A,A'}}{\frac{F_{1,2}}{F_{1,0}} - \frac{F_{2,2}}{F_{2,0}}}$$



Approach 2:

Constructing radial moments  $\delta \langle y_2 \rangle^{A,A'}$  and  $\delta \langle y_4 \rangle^{A,A'}$  that are orthogonal inside the nuclear volume:

$$\delta 
u_{k,\mathrm{FS}}^{(1)\mathcal{A}\mathcal{A}'} = c_{k,0} \delta \langle y_2 \rangle^{\mathcal{A}\mathcal{A}'} + c_{k,2} \delta \langle y_4 \rangle^{\mathcal{A}\mathcal{A}'}$$

•  $c_{k,0}$  and  $c_{k,2}$  are linear combinations of  $F_{k,N}$ •  $\delta \langle y_2 \rangle^{A,A'}$  and  $\delta \langle y_4 \rangle^{A,A'}$  are linear combinations of  $\delta \langle r^2 \rangle^{A,A'}$  and  $\delta \langle r^4 \rangle^{A,A'}$ .



## Simultaneous extraction of $\delta \langle r^2 \rangle$ and $\delta \langle r^4 \rangle$ values



Convergence in frequency field shift as a function of radial moment order using original, but truncated, reformulated field shift (black) and orthogonal radial moments within nuclear volume (purple) in the case of Li-like <sup>200,208</sup>Pb.

## Simultaneous extraction of $\delta \langle r^2 \rangle$ and $\delta \langle r^4 \rangle$ values



Extraction of  $\delta \langle r^4 \rangle$  values from "pseudo-experimental" field shifts in Li-like Pb and Nd isotopes.



- For the  ${}^{2}P_{1/2} {}^{2}S_{1/2}$  and  ${}^{2}P_{3/2} {}^{2}S_{1/2}$  transitions in Li-like Nd  $\delta \nu_{k,\text{FS}}^{A,A'}$  needs to be determined with a relative uncertainty of approximately  $10^{-5}$  in order to extract  $\delta \langle r^{4} \rangle$  with 20% accuracy
- However, using pair of transitions where the electronic factors are less linear dependent a relative uncertainty of  $10^{-4}$  may be sufficient.



25 August 2015

- Knowledge of nuclear shapes and deformation is important for the extraction of accurate  $\delta\langle r^2 \rangle$  values
- $\bullet$  The extraction of  $\delta \langle r^2 \rangle$  values benefits from nuclear Hartree-Fock calculations
- The extraction of  $\delta \langle r^2 \rangle$  and  $\delta \langle r^4 \rangle$  values are in principle possible, but requires more effort on both the experimental and theory side.
- In the future it will be investigated whether the method of reformulated field shift is applicable also to  $K\alpha$  isotope shifts and muonic atoms



# Thank you!



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Nuclear charge and radial electron densities in ground states of selected noble gases





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